Energy Sources in Stars and the Origin of the Elements (15.2-15.3)

What makes a star shine?

Energy Sources in Stars (§15.2)

What makes a star shine?

Sun's Energy Output $L_{sun} = 4 \times 10^{26} \text{ W} (\text{J s}^{-1})$ $t_{sun} = 4.5 \times 10^9 \text{ yr} (1.4 \times 10^{17} \text{ s})$ - oldest rocks - radioactive dating

 $L_{sun} \sim constant over geological timescale (fossil evidence)$

Therefore total energy output $E_{tot} = L_{sun} \times t_{sun} = 6 \times 10^{43} \text{ J}$

How might the Sun have generated this energy?

Energy Source #1:Gravitational contraction
We believe Sun collapsed from a large gas cloud (R ~ ∞) to its present size

Potential energy released is: $\Delta E_{grav} = -(E_{now} - E_{initial})$ take $E_{initial} = 0$ P.E. is obtained from integrating grav. potential over all points in the sphere

P.E. =
$$-\int_{0}^{R_{sun}} G \frac{\left(\frac{4}{3}\pi\rho r^{3}\right)\left(4\pi r^{2}\rho dr\right)}{r}$$

= $-\left(\frac{16}{15}\right)\pi^{2}\rho^{2}GR^{5}$ for constant ρ
With $M = \frac{4}{3}\pi R^{3}\rho$, we get
P.E. = $-\frac{3GM_{sun}^{2}}{5R_{sun}}$



No change in Mass, M_{\odot}

Energy Source #1:Gravitational contraction

 $P.E. = -\frac{3Gm_{sun}^2}{5R_{sun}}$

However, from Virial Theorem ($E_{th} = -1/2 E_{grav}$), only 1/2 of the gravitational energy is available for energy release, the other half heats the star.

So, $\Delta E_{\text{grav}} = -(E_{\text{now}} - E_{\text{initial}})$ = $3/10(GM_{\text{sun}}^2)(1/R_{\text{sun}} - 1/\infty)$ ~ 10^{41} J

 $\sim 1/600 \text{ E}_{\text{tot}}$ needed over Sun's lifetime

By this process, Sun could only radiate for

 $t = \Delta E_{grav} / L_{sun} \sim 10^7$ years

 $t = \Delta E_{grav} / L_{sun}$ is Kelvin-Helmholtz timescale

Conclusion: Gravitational Contraction not the main energy source

in Sun

Energy Source #2:Chemical Reactions

e.g. How much energy would be released if Sun was completely ionized and then all gas recombined into neutral atoms? (binding energy of the H atom 13.6 ev: 1 ev = $1.6 \times 10^{-19} \text{ J}$ $\rightarrow 13.6 \text{ ev} \sim 2 \times 10^{-18} \text{ J}$)

Assume Sun is 100% H (X = 1) Total # H atoms $N_H = M_{sun}/m_H \sim 10^{57}$ $\therefore E_{tot}$ ionization ~ 10^{57} x (2 x 10^{-18}) J ~ 2 x 10^{39} J This is only ~ 1/30,000 E_{total} of Sun's energy output Sun's energy output ~ burning 7000 kg coal each hour on every sq. metre Sun's surface!

Conclusion: Chemical Reactions cannot be source Sun's energy

Energy Source #3:Nuclear Energy

Fission?

Z = 0.02 (2% Sun consists of heavy elements). Fissionable isotopes like ²³⁵U are rare (isotopes are nucleii with the same # protons but different # neutrons) e.g.
n + U²³⁵ → Ba¹⁴¹ + Kr ⁹² + 3n (~3x10⁶ more energy than burning coal)
Fusion?

e.g. 4 protons get converted into 2p + 2n i.e. $4H \rightarrow 1$ ⁴He H is very abundant in the Sun (X = 0.73)

4 H nuclei: mass = $4 \times m_p = 6.693 \times 10^{-27} \text{ kg}$ 1 ⁴He nucleus: mass = $6.645 \times 10^{-27} \text{ kg}$ (binding energy included) Difference is $0.048 \times 10^{-27} \text{ kg} = 0.7\%$ of original mass Where does this mass go?

Energy Source #3:Nuclear Energy

The "lost" mass is converted into energy according to Einstein's equation for the rest energy of matter: $E = mc^2$ (E energy, m mass, c speed light)

Example:

Suppose that Sun was originally 100% H and only 10% of that was available for fusion. Thus...

 $M_{\rm H}$ available = 0.1 $M_{\rm sun}$

 $\Delta M_{\rm H}$ destroyed through fusion = 0.1 M_{sun} x 0.007

total nuclear energy available $E_{tot} = (7 \times 10^{-4} M_{sun})c^2 \sim 1.3 \times 10^{44} J$

Sun's total lifetime energy output = $6 \times 10^{43} \text{ J}$

 $E_{tot} \sim 2 x$ the total energy Sun has emitted!! \rightarrow Sun could shine for at least another 5 x 10⁹ yr at its current luminosity Hence $t_{nuclear}$ for Sun $\sim 10^{10}$ yrs

Protons must overcome mutual electrostatic repulsion: Coulomb

Barrier

Classical Approach:

 E_{kin} proton > $E_{coulomb}$ E_{coulomb} normally written as $k[e_1e_2]/r$ with e's the charge on the protons and k Coulomb's Constant – but if SI units used the charge is defined so that k = 1 and is dimensionless – then e^2/r has units of energy with the value of e (SI) for a proton then $4x10^{-10}$. Hence for a separation between protons of 10^{-13} cm (10^{-15} m) the energy is $(4x10^{-10})^2/10^{-13} =$ $1.6 \times 10^{-6} \text{ ergs} = 1.6 \times 10^{-13} \text{ Joules}$



Classical Approach:

E_{kin} proton > E_{coulomb}
 1/2 m_pv² > e²/r (cgs form)
 LHS is thermal gas energy, thus, 3/2 kT > e²/r

- → At $T_{central} \sim 1.6 \times 10^7 \text{ K} \rightarrow E_{kin}$ ~ 3.3 x 10⁻¹⁶ J
- For successful fusion, r ~ 10^{-15} m (1 fm) \rightarrow

E_{coulomb} ~ 1.6 x 10⁻¹³ J $E_{kin} \sim 10^{-3} E_{coulomb}$ Classically, would need T~10¹⁰ to overcome Coulomb barrier So, no Fusion??



Can we be helped by considering the distribution of energies? - not all protons will have just $E_{th} = 3/2 \text{ kT}$

Proton velocities are distributed according to the Maxwellian equation: $P(v) = 4\pi (m/2\pi kT)^{3/2} v^2 e^{-mv^2/2kT}$

At the high energy tail of the Maxwellian distribution, the relative number of protons, with $E > 10^3 \text{ x } E_{\text{th}}$ is:

 $N(E_{coulomb} \sim 10^3 < E_{thermal} >)/N(< E_{thermal} >) \sim e^{-\Delta E/kT} \sim e^{-1000} = 10^{-430}!!$

Number protons in Sun = $M_{sun}/M_{H} \sim 10^{57}$ so 1 in 10^{430} of these will have enough energy to overcome Coulomb barrier.

So again, no nuclear reactions??

Quantum Approach:

Heisenberg Uncertainty Principle - $\Delta p \Delta x > h/2\pi$

If Δp small, then Δx may be large enough that protons have non-negligible probability of being located within 1 fm of another proton inside Coulomb barrier

Called Quantum Tunneling

- Particles have wavelength, $(\lambda = h/p)$, associated with them (like photons) called de Broglie wavelength
- Proved in many experiments for example diffraction electrons (wave phenomenon)
- Eg. λ free electron (3 x 10⁶ m/s) = 0.242 nm (size atom)
- λ person (70 k gm) jogging at 3 m/s = 3 x 10⁻³⁶ m (negligible - person won' t diffract!)
- Increased wavelength reflects loss momentum.

Classically, the proton is reflected at R_0 .

Quantum mechanically, there exists a finite probability for a proton to reach interaction radius, R.



Even with tunneling, are stars hot enough for nuclear reactions to proceed?

 $\lambda = h/p$ is the wavelength associated with a massive particle (p = momentum)

In terms momentum, the kinetic energy of a proton is: $1/2 m_p v^2 = p^2/2m_p$

Assuming proton must be within 1 de Broglie λ of its target to tunnel set distance, r, of closest approach to λ , (where barrier height = original K.E. incoming particle) gives:

 $e^{2}/r = e^{2}/\lambda = p^{2}/2m_{p} = (h/\lambda)^{2}/2m_{p}$

Solving for λ (=h²/2m_pe²) and substituting r = λ into 3/2 kT = e²/r, we get the QM estimate of the temperature required for a nuclear reaction to occur: T_{QM} = 4/3(e⁴m_p/kh²) - putting in numbers

 $T_{QM} = 10^7$ K which is comparable to T_{core} . Therefore, fusion is feasible at centre of Sun



Basic particles involved in nuclear reactions

Particle	Symbol	Rest mass (kg)	Charge (e-)	Spin
photon	γ	0	0	1
neutrino	ν	>0	0	1/2
anti-neutrino	V-	>0	0	1/2
electron	e-	9 x 10 ⁻³¹	-1	1/2
positron	e+	9 x 10 ⁻³¹	+1	1/2
proton $(^{1}_{1}H)$	p+	1.6 x 10 ⁻²⁷	+1	1/2
neutron	n	1.6 x 10 ⁻²⁷	0	1/2
deuteron	$^{2}_{1}H$	3.2 x 10 ⁻²⁷	+1	

Conservation Laws in nuclear reactions:

(1) mass-energy

(2) charge

(3) difference between number particles and anti-particles conserved ie particle cannot be created from anti-particle or vice-versa but a pair can be formed or destroyed without violating this rule or #p + #n conserved



Types of nuclear reactions

Beta Decay: $n \rightarrow p^+ + e^- + v^-$ proceeds spontaneously - also for neutron inside nucleus eg [Z-1, A] \rightarrow [Z, A] + $e^- + v^-$ (Z = # p⁺, N = # n, A = Z+N) ${}^{14}_{6}C \rightarrow {}^{14}_{7}N + e^- + v^-$

Inverse Beta Decay: $p^+ + e^- \rightarrow n + \nu$ eg [²⁶₁₃Al + e⁺ \rightarrow ²⁶₁₂Mg + ν]

(p⁺, γ) process: ^AZ + p⁺ \rightarrow ^{A+1}(Z+1) + γ eg [¹²C + p⁺ \rightarrow ¹³N + γ]

(α , γ) process: α particle (⁴He) added to nucleus to make heavier particle [^AZ + ⁴He \rightarrow ^{A+4}(Z+2) + γ] eg. [⁸Be + ⁴He \rightarrow ¹²C + γ]

Since the Sun's mass consists mostly of H and He, we anticipate nuclear reactions involving these two elements eg $p^+ + p^+ \rightarrow {}^{2}\text{He}$ $p^+ + {}^{4}\text{He} \rightarrow {}^{5}\text{Li}$ ${}^{4}\text{He} + {}^{4}\text{He} \rightarrow {}^{8}\text{Be}$

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But there is a problem here - what is it? All these reactions produce unstable particles eg $p^+ + p^+ \rightarrow {}^{2}\text{He} \rightarrow p^+ + p^+$ $p^+ + {}^{4}\text{He} \rightarrow {}^{5}\text{Li} \rightarrow p^+ + {}^{4}\text{He}$ ${}^{4}\text{He} + {}^{4}\text{He} \rightarrow {}^{8}\text{Be} \rightarrow {}^{4}\text{He} + {}^{4}\text{He}$

So among lightest elements there are no two-particle exothermic reactions producing stable particles. We have to look to more peculiar reactions or those involving rarer particles

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Proton-Proton Chain

Using considerations of: — minimizing Coulomb barriers,

- crossections,
- and making sure conservation laws are obeyed
- the nuclear reaction chain at right produces the energy observed in the Sun.



Hans Bethe 1967

The proton-proton chain PPI Branch





Summary of proton-proton chain:

- $6p^+ \rightarrow {}^{4}\text{He} + 2p^+ + 2e^+ + 2\nu + 2\gamma \text{ or } 4p^+ \rightarrow {}^{4}\text{He} + 2e^+ + 2\nu + 2\gamma$
- Mass difference between 4p⁺ and 1 ⁴He = 26.7 Mev / 6.242 x 10¹⁸ ev/J = $4.2 \times 10^{-12} \text{ J}$
- ~3% of this energy (0.8 Mev) is carried off by neutrinos and does not contribute to the Sun's luminosity
- 2 e⁺ immediately annihilate with 2 e⁻ and add 4 x 0.511 Mev (= 2.04 Mev)
- So, the total energy available for the Sun's luminosity per ⁴He formed is: (26.7 0.8 +2.04) Mev = 27.9 Mev = 4.3 x 10⁻¹² J
- # ⁴He formed/sec = $L_{sun}/4.3 \times 10^{-12}J = 3.9 \times 10^{26} J/s / 4.3 \times 10^{-12}$ J = 9.0 x 10^{37 4}He/s
- Increase of ⁴He mass/time = $dm_{He}/dt = 9.0 \times 10^{37} {}^{4}\text{He/s} \times 6.68 \times 10^{-27} \text{ kg}/{}^{4}\text{He} = 6.0 \times 10^{11} \text{ kg} {}^{4}\text{He}/\text{s}$
- After 10^{10} years (3 x 10^{17} sec), M(⁴He) = 1.8 x 10^{29} kg = 10% mass Sun

Check on Proton-Proton Chain

- # ⁴He formed/sec = 9.0 x 10³⁷ ⁴He/s and each ⁴He is accompanied by production of 2 v → 2.0 x 10³⁸ v /s
- Flux neutrinos at Earth is 2 x $10^{38}/4\pi R_{ES}^2 = 7 \times 10^{14} \text{ v m}^{-2} \text{ s}^{-1}$
- v can be capture at Earth by $v + {}^{37}Cl \rightarrow {}^{37}Ar + e$
- Prediction was 5.8/day detected 2.1/day
- Concern that T Sun was thus lower than model



 Answer was 3 flavors of v change one to other inside Sun – new physics – finally resolved in SNO which was not sensitive to v types – detected expected # of v. Ray Davis 2002 Art McDonald 2015



The previous nuclear reaction chain (PPI) is not the only way to convert H into He.

eg last step: ³He + ³He \rightarrow ⁴He + 2p⁺ can proceed differently if there is appreciable ⁴He present

A second branch (PPII)

In solar centre, 69% of ${}_{2}^{3}$ He combines with another ${}_{2}^{3}$ He, but 31% can fuse with ${}_{2}^{4}$ He

 ${}_{2}^{3}\text{He} + {}_{2}^{4}\text{He} \rightarrow {}_{4}^{7}\text{Be} + \gamma$ ${}_{4}^{7}\text{Be} + e^{-} \rightarrow {}_{3}^{7}\text{Li} + v_{e}$ ${}_{3}^{7}\text{Li} + p^{+} \rightarrow {}_{2}^{4}\text{He} + {}_{2}^{4}\text{He}$ <u>A</u> third branch (PPIII) In the Sun's core, a ${}_{4}^{7}$ Be nucleus can capture a p⁺ instead of an e⁻ about 0.3% of the time. $\sum_{A}^{7} Be + p^{+} \rightarrow {}_{5}^{8} B + \gamma$ Very unstable $\rightarrow {}_{5}^{8}B \rightarrow {}_{4}^{8}Be + e^{+} + v_{e}$ ${}^{8}_{4}\text{Be} \rightarrow {}^{4}_{2}\text{He} + {}^{4}_{2}\text{He}$ **PPI + PPII + PPIII operate** simultaneously





Alternate Fusion Reactions $H \rightarrow {}^{4}He$

CNO Cycle (or bi-cycle)

- First step in PP chain has very low reaction rate (8x10⁹ years - weak interaction).
- ¹²C can act as a catalyst in fusion reaction.
- Net result:

 $^{12}C + 4p^+ \rightarrow ^{12}C + ^{4}He + 2e^+ + 2v + 3\gamma$

- Note: ¹²C neither created nor destroyed. Also isotopes of N & O are temporarily produced.
- CNO cycle dominates over PP chain if $T_{core} > 1.8 \ge 10^7$ (slightly hotter than Sun)

THE CNO BI-CYCLE				
(T < 10 ⁸ K)				
$\mathbf{f} \mathbf{E}^{12} + \mathbf{H}^{1} \rightarrow \mathbf{N}^{13} + \gamma$	10 ⁶ yr			
$N^{13} \rightarrow C^{13} + e^+ + \nu$	14 min			
$C^{13} + H^1 \rightarrow N^{14} + \gamma$	3 x 10⁵ yr			
$ N^{14} + H^1 \rightarrow O^{15} + \gamma $	3 x 10 ⁸ yr			
$O^{15} \rightarrow N^{15} + e^+ + \nu$	82 s			
$N^{15} + H^1 \rightarrow C^{12} + He^4$	10 ⁴ yr			
or (~4 × 10 ⁻⁴)	X as frequently			
$N^{15} + H^1 \rightarrow O^{16} + \gamma$				
$O^{16} + H^1 \rightarrow F^{17} + \gamma$	Once for every 2500			
$F^{17} \rightarrow O^{17} + e^+ + \nu$	instances o			
$O^{17} + H^1 \rightarrow N^{14} + He^4$	Cycle 1			



Helium Burning

No stable nuclei with A=8 therefore He burning proceeds in 2 steps

1.⁴He + ⁴He $\leftarrow \rightarrow$ ⁸Be – builds up small concentration ⁸Be 2.⁸Be + ⁴He \rightarrow ¹²C – possible at T=10⁸ K

Net effect is ³He4 \rightarrow ¹²C + γ

When sufficient ¹²C is produced we can have ${}^{12}C + {}^{4}He \rightarrow {}^{16}O$

Carbon Burning and Beyond $^{12}C + ^{12}C \rightarrow ^{24}Mg + \gamma \text{ or }^{20}Ne + ^{4}He \text{ or }^{23}Na + p$

Can be followed by ${}^{20}\text{Ne} + \gamma \rightarrow {}^{16}\text{O} + {}^{4}\text{He or } {}^{20}\text{Ne} + {}^{4}\text{He} \rightarrow {}^{24}\text{Mg}$

Oxygen burning ${}^{16}O + {}^{16}O \rightarrow {}^{32}S + \gamma \text{ or } {}^{28}Si + {}^{4}He \text{ or } {}^{31}P + p \text{ (many other channels)}$

Si burning does not occur directly but in a complex way by photodisintegrations (alpha, γ) and alpha capture \rightarrow core ⁵⁶Fe as it is most tightly bound of all the nuclei eg ²⁸Si + ⁴He \rightarrow ³²S, ³²S + ⁴He \rightarrow ³⁶Ar, ³⁶Ar + ⁴He \rightarrow ⁴⁰Ca, ⁴⁰Ca + ⁴He \rightarrow ⁴⁴Ti, ⁴⁴Ti + ⁴He \rightarrow ⁴⁸Cr, ⁴⁸Cr + ⁴He \rightarrow ⁵²Fe ⁵²Fe + ⁴He \rightarrow ⁵⁶Ni (these kinds of reactions produce many of the heavy elements up to Fe that we see in the Universe)



Summary Other Fusion Reactions

Fusion of	into	at T _{central}	Ref.
⁴ He triple α process	Be, C	~10 ⁸ K	p. 368
¹² C	O, Ne, Na, Mg	6 x 10 ⁸ K	p. 409
¹⁶ O	Mg, Si, S, etc.	~10 ⁹ K	p. 409

We believe Universe began with only H, He, (Li, Be) All other elements created in core of stars (stellar nucleosynthesis) We are all made of "STARSTUFF"

Elements beyond iron: s-process

Where do the elements beyond Fe come from (eg gold = ${}^{197}Au_{79}$?) Coulomb barriers too high for expected stellar central temperatures for reactions like ${}^{56}Fe + {}^{56}Fe \rightarrow$?



Elements beyond iron: s-process

Where do the elements beyond Fe come from (eg gold = ${}^{197}Au_{79}$?) Coulomb barriers too high for expected stellar central temperatures for reactions like ${}^{56}Fe + {}^{56}Fe \rightarrow$?

Solved by understanding of the s-process (slow neutron capture) where a heavy nucleus captures a neutron to make a heavier isotope. If the isotope is stable, it can capture another neutron until the nucleus becomes unstable at which point it can β -decay to produce an element of the next highest atomic number. These are NOT charged particle reactions so problem of Coulomb barriers is not an issue. Fe is the seed of the s-process and a source of neutrons is required. The stars wherein this process occurs are luminous giant stars.

Elements beyond iron: s-process

Neutron sources: ${}^{13}C + {}^{4}He \rightarrow {}^{16}O + n \text{ or } {}^{22}Ne + {}^{4}He \rightarrow {}^{25}Mg + n$

s-process example: ⁵⁶Fe + n \rightarrow ⁵⁷Fe, capture 2 more n \rightarrow ⁵⁹Fe, Fe 56-58 are all stable but ⁵⁹Fe unstable (half life 44.5 days).

If flux n not too high, and don't capture another n within 44.5 days, ⁵⁹Fe can decay to ⁵⁹Co by β -decay (n \rightarrow p + e⁻ + v) and we now have an element one element up in periodic table.

The process can continue to form higher and higher mass elements and thus we can populate the high mass end of the periodic table.

Elements beyond iron: r-process

In contrast to the scenario just described, imagine that the flux of neutrons is high and the time between neutron captures is small compared to the half life of the isotopes under consideration, super-rich isotopes can then be formed. When the neutron flux stops, these super-neutron-rich isotopes will undergo a series of β -decays until a stable isotope is reached. As you might imagine, this scenario can take place during a supernova explosion (more about this soon).

For example, here is a scenario to make GOLD

¹⁹⁵Y₇₀ + n \rightarrow ¹⁹⁶Y₇₀, ¹⁹⁶Y is unstable and decays ¹⁹⁶Y \rightarrow ¹⁹⁶Lu₇₁ + e⁻ + v Lu (Lutecium) then captures a n to make ¹⁹⁷Lu which is unstable and it decays to ¹⁹⁷Hf₇₂. After multiple n captures and decays we end up with ¹⁹⁷Pl₇₈ which then β -decays ¹⁹⁷Pl₇₈ \rightarrow ¹⁹⁷Au₇₉ + e⁻ + v. ¹⁹⁷Au₇₉ is the one stable isotope of GOLD.

Elements beyond iron: r-process: Collision Neutron Stars as Neutron Source



Origin of the Elements

