

APPENDIX A

Derivation of the Detection Sensitivity of an Astronomical Telescope

We first consider the completely general case for an arbitrary telescope, detector, and observing conditions. Next we consider the NASA-LMT in two specific operating modes as they relate to astronomy and orbital object detection.

Signal:

From first principles, the total signal electrons (S) registering on a detector from an external source is given by:

$$S = A_{TEL} \cdot E_S \int_{BW} TE(\lambda) \cdot QE(\lambda) \cdot F_S(\lambda) d\lambda \quad (AA.a)$$

Where:

A_{TEL} = Telescope primary aperture

E_S = Signal exposure time

BW = Detector bandwidth

$TE(\lambda)$ = Telescope efficiency or throughput as a function of wavelength

$QE(\lambda)$ = Detector quantum efficiency as a function of wavelength

$F_S(\lambda)$ = Signal flux as a function of wavelength

λ = Wavelength

For a point source we can rewrite $F_s(\lambda)$ in terms of apparent magnitude $m(\lambda)$:

$$F_s(\lambda) = N \cdot 10^{(-0.4 \cdot m(\lambda))} \quad (\text{AA.b})$$

Where:

$N=10,000 \text{ } \gamma/\text{sec cm}^2\text{nm}$ (the photon flux for 0th apparent magnitude)

For a diffuse source we can rewrite $F_s(\lambda)$ in terms of surface brightness $m_{SB}(\lambda)$:

$$F_s(\lambda) = N \cdot A_{PIX} 10^{(-0.4 \cdot m_{SB}(\lambda))} \quad (\text{AA.c})$$

Where:

A_{PIX} = Area of the detector or detector element (arc second^2)

Making the assumption that the source flux and telescope and detector response are relatively independent of wavelength over each B,V,R, or I Band optical wavelength regime, we can approximate A-1 as:

$$S = A_{TEL} \cdot BW \cdot TE \cdot QE \cdot E_s \cdot F_s \quad (\text{AA.d})$$

Expressing in terms of the actual signal flux at the detector S_F :

$$S = E_S \cdot S_F \quad (\text{AA.e})$$

Where:

$$S_F = A_{TEL} \cdot BW \cdot TE \cdot QE \cdot F_S \quad (\text{AA.f})$$

If this signal is spread uniformly over a n pixel region of the detector in units of S_{pix} ,

we can write:

$$S = n \cdot S_{pix} = n \cdot E_{S_{pix}} \cdot S_{F_{pix}} \quad (\text{AA.g})$$

Where:

$$E_{S_{pix}} \equiv \frac{E_S}{\alpha} = \text{Exposure time per pixel}$$

$$S_{F_{pix}} \equiv \frac{S_F}{\beta} = \text{Source flux per pixel}$$

α = exposure reduction parameter due to object motion

β = signal flux reduction parameter due to image spread

With the constraint $n = \alpha \cdot \beta$ and $\alpha, \beta \geq 1$

Depending on the nature of the source signal, the per pixel components of signal and exposure time can vary. Equation AA.g accounts for any configuration of source and detector: static point source sampled by a single pixel, static diffuse source sampled by multiple pixels, moving source sampled by single pixels, moving diffuse source sampled by multiple pixels.

Noise:

Governed by Poisson photon counting statistics, the total noise associated with a signal detection by a CCD detector is given by:

$$N_T = \sqrt{(R_T^2 + S + BKG + D_T)} \quad (\text{AA.h})$$

Where:

R_T = Total CCD read noise for all signal pixels (includes Bias noise)

S = Source signal

BKG = Total Background

D_T = Total Thermal Dark charge

If the signal is spread evenly over n pixels, and the CCD response is uniform, we can rewrite the total noise as:

$$N_T = \sqrt{n \cdot (R_{pix}^2 + S_{pix} + BKG_{pix} + D_{pix})} \quad (AA.i)$$

Where:

R_{pix} = CCD read noise per pixel

S_{pix} = Source signal per pixel

BKG_{pix} = Background per pixel

D_{pix} = Thermal Dark charge per pixel

The amount of time the detector is exposed to the field is not necessarily equal to the amount of time the signal source is exposed to the detector (eg. a meteor transit during a sidereal drift scan). Rewriting the noise in terms of the total detector exposure E_D and signal exposure per pixel $E_{S_{pix}}$:

$$N_T = \sqrt{n \cdot (R_{pix}^2 + E_{S_{pix}} \cdot S_{F_{pix}} + E_D \cdot BKG_{F_{pix}} + E_D \cdot D_{C_{pix}})} \quad (AA.j)$$

Where:

R_{pix} = CCD read noise per pixel

E_D = Detector exposure time

$BKG_{F_{pix}}$ = Background flux per pixel

$D_{c_{pix}}$ = Thermal Dark current per pixel

Signal-to-Noise Ratio:

The total signal-to-noise-ratio (SNR) can be expressed as the quotient of equations

(AA.g) and (AA.j):

$$SNR = \frac{n \cdot E_{S_{pix}} \cdot S_{F_{pix}}}{\sqrt{n \cdot (R_{pix}^2 + E_{S_{pix}} \cdot S_{F_{pix}} + E_D \cdot BKG_{F_{pix}} + E_D \cdot D_{c_{pix}})}} \quad (AA.k)$$

Since we are interested in detection limits where the signal is small, we may safely

neglect the signal related component of the noise. ($S_{pix} \ll R_{pix}^2, BKG_{pix}$)

Simplifying, we have:

$$SNR = \frac{\sqrt{n} \cdot E_{S_{pix}} \cdot S_{F_{pix}}}{\sqrt{(R_{pix}^2 + E_D \cdot BKG_{F_{pix}} + E_D \cdot D_{C_{pix}})}} \quad (AA.1)$$

Examples:

Let us consider the case of the detection sensitivity, in two operating modes of the 3 m NASA-LMT with the LSP 2K detector, MgF_2 coated camera head window, and the original 3 element MgF_2 coated corrector lens.

For all cases, the NASA-LMT and LSP 2K CCD have the following operating parameters:

Primary mirror area = $A_{TEL} = 70685.8 \text{ cm}^2$ (neglecting PF obscuration of 2918 cm^2)

CCD Array size = 2048x2048 pixels

Pixel size = 15 μm

Read Noise = $R_{pix} = 28.1 \text{ e-}$

Dark Current = $D_{C_{pix}} = 0.08 \text{ e-/pix-sec}$

Telescope efficiency = $TE = 0.679$

(76.6% Hg reflectivity @ V and eight MgF_2 surfaces @ 98.5% each)

$QE = \sim 0.35$ (V Band; 500-600 nm); ~ 0.18 (White Light; 450-975 nm)

$BW = 100 \text{ nm}$ (V Band; 500-600 nm); 525 nm (White Light; 450-975 nm)

Case 1)

Find the limiting stellar magnitude (V Band) for a sidereal drift scan with no CCD binning (1x1).

We wish to solve equation (AA.1) for $S_{F_{pix}}$ and then use equations (AA.b) and (AA.f)

to find the corresponding limiting magnitude.

The exposure time for a zenith sidereal drift-scan is determined by the telescope latitude, CCD array size, and plate scale.

First, the Sidereal Drift Rate (SDR) is given by:

$$\text{SDR (arcseconds/sec)} = (2\pi / 86,160) \cdot \cos \delta \cdot 206,265 \quad (\text{AA.m})$$

Where 86,160 is the number of seconds per sidereal day and 206,265 is the number of arcseconds/radian.

At the NODO location, with latitude $\delta = 33$ degrees North, this yields:

$$\text{SDR} = 12.615 \text{ arcseconds/sec.}$$

Second, the plate scale (PS) is determined by the CCD pixel size and the telescope focal length:

$$PS \text{ (arcseconds/pixel)} = 206,265 \cdot \left(\frac{PixelSize}{FocalLength} \right) \quad (AA.n)$$

For the NASA-LMT with an effective (after the finite power corrective optics) focal length of 5,173.87 mm (f/1.7246) and a LSP 2K CCD pixel size of 15 μ m, we have:

$$PS = 0.598 \text{ arcseconds/pixel}$$

Finally, the field crossing time (FCT), and hence the total exposure time, for a star transiting the FOV is given by:

$$FCT = \frac{CCDArraySize \cdot PlateScale}{SiderealDriftRate} \quad (AA.o)$$

For the NASA-LMT and LSP 2K CCD, this yields:

$$FCT = 97.1 \text{ sec}$$

Since we are drift scanning at the sidereal rate, the stellar images are static relative to the electro-optically tracking detector pixels ($\alpha=1$), therefore:

$$FCT = E_S = E_D = E_{S_{pix}} = 97.1 \text{ sec}$$

The per pixel background flux $BKG_{F_{pix}}$ is dependent on the telescope location.

For the NODO location, the night sky background is approximately (Schneeberger et al. 1979):

$$\text{NODO Sky BKG (V band)} = 21.81 \text{ magnitudes / arc second}^2$$

Substituting into equation (AA.c) and safely neglecting any wavelength dependence

within the V band, we have (since $A_{PIX} = PS^2 = 0.36 \text{ arc second}^2$):

$$F_S(\lambda) = F_S = 6.797 \times 10^{-6} \gamma / \text{sec cm}^2 \text{ nm}$$

Substituting into equation (AA.f), the corresponding background flux at a detector pixel is:

$$BKG_{F_{pix}} = 11.42 \text{ e-/pix-sec}$$

Substituting all known values into equation (AA.l):

$$SNR = 2.229 \sqrt{n} \cdot S_{F_{pix}}$$

The nominal NODO seeing value is approximately 1.2 arcseconds FWHM which implies that a stellar Point Spread function (PSF) could occupy as few as 4 pixels ($n = \beta = 4$). If we use the generally accepted limiting SNR of 6.3 we find:

$$S_{F_{pix}} = 1.413 \text{ e-/pix} \Rightarrow S_F = 5.653 \text{ e-}$$

Using equation (AAf), this corresponds to a source flux F_S of:

$$F_S = 3.365 \times 10^{-6} \text{ } \gamma / \text{sec cm}^2 \text{ nm}$$

We use equation (AA.b) to recover the limiting V apparent magnitude:

$$m_V = 23.68$$

At a SNR of 10 and a PSF spread over 9 pixels ($n = \beta = 9$), corresponding to 1.8 arcsecond seeing, we obtain:

$$m_V = 22.74$$

V Band observations with the NASA-LMT are made through a glass filter with a bandwidth of approximately 100 nm and an average transmittance of 0.35.

This yields actual limiting magnitudes of:

$$m_V = 22.54 \text{ for } 1.2 \text{ arcsecond seeing}$$

$$m_V = 21.60 \text{ for } 1.8 \text{ arcsecond seeing}$$

These values correlate well with empirical results, some of which are given in Chapter VI.

Case 2)

Find the limiting magnitude (V Band) and detection diameter for a high-speed drift-scan of an orbital object at 800 km altitude.

Physical limitations intrinsic to the LSP 2K CCD detector require using binned pixels in order to implement a high-speed drift-scan. Binning reduces the number of pixels which must be read by combining the charge from adjacent rows (parallel binning) or columns (serial binning) into super-pixels. Since the readout electronics require a finite amount of time to digitize each pixel, the CCD surface can be read more quickly if there are fewer pixels. Excessive binning has the disadvantage, however, of increasing the background noise associated with each portion of signal. From a SNR perspective, the signal PSF and pixel (or super-pixel) size should match, although from a photometric perspective, the signal PSF should be over-sampled by a factor of two. This reduces SNR slightly, but increase photometric accuracy by reducing the effect of pixel-to-pixel sensitivity variations.

Through extensive testing, the timing equation for the LSP 2K CCD was found empirically to be given by:

$$\text{SFRT} = \text{Single Frame Read Time (sec)} = \frac{4 \cdot \text{Delay} + 10.2}{\text{Bin}} \quad (\text{AA.p})$$

Where:

Delay = parallel shift delay time (sec)

Bin = symmetrical binning factor in which:

$$Bin = \text{serial binning}(Bin_s) = \text{parallel binning}(Bin_p)$$

To optimize high-speed drift-scan object tracking, we require that the SFRT equal the object FCT. The FCT is given by:

$$FCT = \frac{CCDArraySize \cdot PlateScale}{\varpi_{EFF}} \quad (AA.q)$$

Where:

ϖ_{EFF} = effective orbital angular velocity of object accounting for the earth's rotation

Assuming a circular orbit, the angular velocity ϖ_{EFF} (radians/sec) is given by the following expressions derived from simple gravitational and centrifugal force balance and allowing for the east-west motion of the telescope due to the earth's rotation:

$$\varpi_{EFF}^2 = \varpi_{NS}^2 + \varpi_{EW}^2 \quad (AA.r)$$

Where:

$$\varpi_{NS} = \varpi_{CIRC} \cdot \sin(90 - PA) \quad (\text{north-south angular rate})$$

$$\omega_{EW} = [(\omega_{CIRC} \cdot \cos(90 - PA)) - \omega_{Earth} \cdot \cos(LAT)] \quad (\text{east-west angular rate})$$

and

$$\omega_{CIRC} = \sqrt{\frac{GM}{H^2(H + R_E)}} \quad (\text{angular rate for a circular orbit}) \quad (\text{AA.s})$$

and

$$\omega_{Earth} = \text{angular rotational velocity of the earth} = 0.004178 \text{ deg/sec}$$

$$PA = \text{position angle of object motion} \quad (\text{north} = 0 \text{ degrees})$$

$$LAT = \text{latitude of the telescope} \quad (\text{NODO} = 32.979 \text{ degrees N})$$

$$G = \text{Gravitational constant} = 6.672 \times 10^{-11} m^3 / kg \cdot sec^2$$

$$M = \text{Mass of the Earth} = 5.98 \times 10^{24} kg$$

$$H = \text{Orbital altitude}$$

$$R_E = \text{Radius of the Earth} = 6378 \text{ km}$$

Solving for ω_{EFF} for an object at 800 km altitude and 0 degree position angle (PA):

$$\omega_{EFF} = 0.6388 \text{ degrees/sec}$$

Substituting into equation(AA.q) yields:

$$FCT = 0.534 \text{ sec}$$

Since we are drift scanning at the orbital object angular rate, the stellar images are static relative to the electro-optically tracking detector pixels ($\alpha=1$), therefore:

$$\text{FCT} = E_S = E_D = E_{S_{pix}} = 0.534 \text{ sec}$$

Equating FCT to SFRT and solving equation (AA.q) for the requisite minimum binning factor, we have:

$$\text{Bin} = 16 \text{ with Delay} = 0.$$

The binning factor must be accounted in the SNR relationship (AA.l). By simply translating the per pixel definitions to per super-pixel definitions, we account for arbitrary binning:

$$\text{SNR} = \frac{\sqrt{n} \cdot E_{S_{spix}} \cdot S_{F_{spix}}}{\sqrt{(R_{spix}^2 + E_D \cdot BKG_{F_{spix}} + E_D \cdot D_{C_{spix}})}} \quad (\text{AA.t})$$

Where:

n = number of super-pixels occupied by the signal

$E_{S_{spix}}$ = signal exposure per super-pixel

$S_{F_{spix}}$ = signal flux per super-pixel

$$R_{spix} = \text{read noise per super-pixel} (= R_{pix})$$

$$BKG_{F_{spix}} = Bin_S \cdot Bin_P \cdot BKG_{F_{pix}} = \text{background level per super-pixel}$$

$$D_{C_{spix}} = Bin_S \cdot Bin_P \cdot D_{C_{pix}} = \text{dark current per super-pixel}$$

As in the first example, we substitute the known quantities below into equation (AA.t):

$$FCT = E_S = E_D = E_{S_{pix}} = 0.534 \text{ sec} = E_{S_{spix}}$$

$$BKG_{F_{pix}} = 11.42 \text{ e-/pix-sec (V Band)}$$

Substituting yields:

$$SNR = 0.01099 \sqrt{n} \cdot S_{F_{pix}}$$

The nominal NODO seeing value is approximately 1.2 arcseconds FWHM which implies that a stellar Point Spread function (PSF) could occupy as few as 4 un-binned pixels. This implies that the orbital object PSF should occupy only 1/16 of the area of a single 16 binned super-pixel, therefore ($n = \beta = 1$). If we use the generally accepted limiting SNR of 6.3 we find:

$$S_{F_{pix}} = 573.3 \text{ e-/pix} = S_F$$

Using equation (AA.f), this corresponds to a V Band source flux F_S of:

$$F_S = 3.413 \times 10^{-4} \text{ } \gamma / \text{sec cm}^2 \text{ nm}$$

We use equation (AA.b) to recover the limiting V apparent magnitude:

$$m_V = 18.67$$

Next let us determine the diameter on an orbital object at 800 km altitude with an apparent magnitude of $m_V = 18.67$.

Since a telescope (as opposed to a RADAR), observes orbital debris via reflected sunlight, the flux received from the object ($F_{s_{OBJ}}$) depends on the solar flux (F_{SUN}) impinging on the object:

$$F_{s_{OBJ}} = F_{SUN} \cdot \frac{A \cdot (\pi \cdot (D/2)^2)}{H^2} \cdot PF \quad (\text{AA.u})$$

Where:

A = Bond albedo (as opposed to geometric albedo)

D = Object diameter

H = Object altitude

PF = Object phase function

Rewriting in terms of apparent magnitudes and solving for D:

$$D = \frac{2H}{\sqrt{\pi \cdot A \cdot PF}} \cdot 10^{(-0.2 \cdot (m_{OBJ} - m_{SUN}))} \quad (AA.v)$$

Where:

$$m_{SUN} = -26.74 \text{ (V Band)}$$

For this example we will use the mean optical albedo $\langle A \rangle$ of 0.1 derived from extensive optical observations of orbital debris objects and comparison with catalogued radar cross-sections. (Henize et al. 1993)

Generally three types of phase function (PF) are considered: the specular, the diffusely reflecting (Lambert) sphere, and the lunar phase function. Each function is normalized to unity over all solid angle.

The specular sphere radiates equally into 4π steradians independent of phase angle (ϕ_p = angular separation between the source and the observer with the object at the vertex) and thus its phase function is simply:

$$PF_S = \frac{1}{4\pi} \quad (AA.w)$$

The Lambert phase function is a diffuse radiator with a maximum in the backscatter direction ($\phi_p = 0$) and a minimum in the forward scattering direction ($\phi_p = \pi$). The

functional relationship is given by:

$$PF_L = \frac{2}{3\pi^2} \left[\sin(\phi_P) + (\pi - \phi_P) \cdot \cos(\phi_P) \right] \quad (\text{AA.x})$$

The lunar phase function is empirically derived from lunar observations (Allen 1973). It has a very strong backscatter component ($\phi_P=0$) and can be approximated in the range ($40 < \phi_P < 110$ degrees) as (Mulrooney 1993):

$$PF_M = \frac{5}{\pi^2} \left[5.3095 \times 10^{-5} \cdot (\phi_P)^2 - 1.2591 \times 10^{-2} \cdot \phi_P + 0.79121 \right] \quad (\text{AA.y})$$

The phase function for orbital objects can generally be approximated by linear combinations of these three functions with most objects characterized by the specular and Lambert functions.

Figure AA-1 illustrates the functional behavior of these three phase functions as well as that of the lunar quadratic approximation.

Since most orbital object observations are made at solar phase angles between 40 and 100 degrees, we consider for this example a Lambert sphere at $\phi_P = \frac{\pi}{2}$, therefore:

$$PF_L = 0.0675.$$

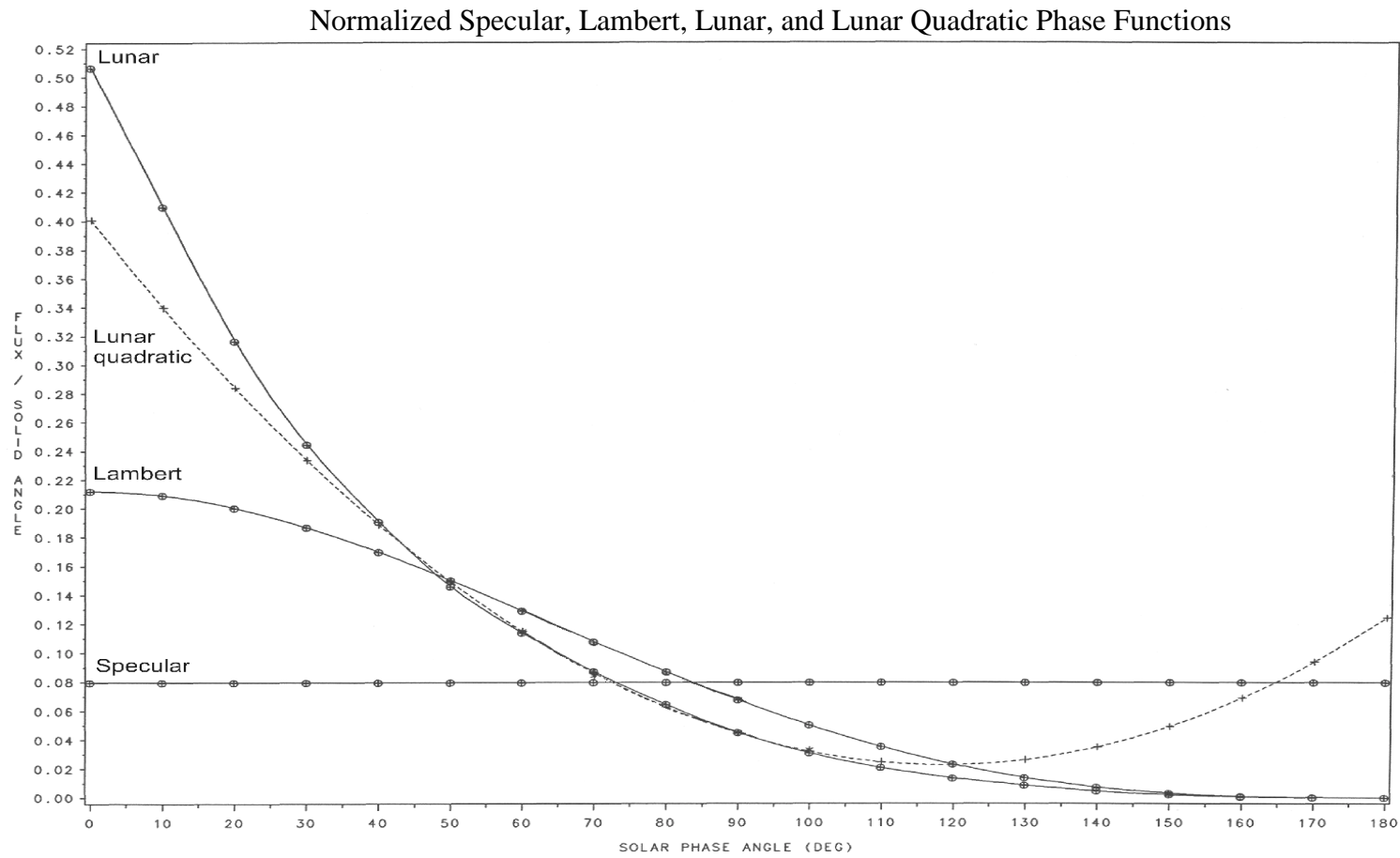


Figure AA-1. A graph demonstrating the three primary phase functions employed in orbital object research. (All are normalized to unity.) The strong backscatter behavior of the lunar function is evident as is the isotropic behavior of the specular sphere. The quadratic approximation to the lunar function fits well in the critical (solar phase angle= 40-100 deg) regime where most optical observations are made. The phase functions of most orbital objects can be approximated by a linear combination of the specular and Lambert functions. (Mulrooney 1993)

Substituting all known values into equation (AA.v) and solving for D:

$$D = 0.910 \text{ cm}$$

Thus ideally, the NASA-LMT can detect an orbital object of 0.910 cm diameter while operating in the drift-matching (DM) mode in the V Band. In reality, this limiting diameter is not achieved, due to the additional noise contribution from background stars on the rapidly drift-scanning pixel. There is also a smearing effect which causes the signal photons to spread over several super-pixels with a consequent reduction in the SNR. This smearing is due to imperfect charge- transfer-efficiency (CTE) in the LSP 2K CCD, the discrete nature of the CCD parallel shifting relative to the continuous motion of the orbital object (or star) (Gibson and Hickson 1992b), and the inability to accurately set the CCD read rate to exactly match the angular velocity of high-speed motion. These limitations normally result in the signal occupying a minimum of 9 super-pixels with a significant loss in sensitivity. For an orbital object with identical properties to those in this example observed with the LSP 2K CCD, the empirical limiting diameter in DM mode is approximately 2.5 cm (Chapter II, Figure II.B-7).

APPENDIX B

(NASA-LMT Orbital Object Detection in Drift Matching (DM) Mode)

```

PROGRAM LMTTDI  FORTRAN 77
! MARK MULROONEY V2.5

CHARACTER*250 DCHR(100)

INTEGER  OUTFILE
REAL  DARK, DR, R, A, B, M2, EXPHLD, LRO
REAL  BIN, RE, GM, PS, FW, PXA, SN, READ, SKY, PF, ALB
REAL  NPIX, WS, WO, INC, ALPHA, RS, RO, RCM, D, PI, DU, DD
REAL  SIGEXP, SOLMAG, NOISE, DELTR, DELTI, M1
REAL  WOT, WEW, WNS, WEARTH, SITE, WST, WEWT, WNST, SI
REAL  SIGEXPMX, EPSILON, EXPMX
REAL      DATA(40,40)
INTEGER  INCMX, RANMX, J1, INCMX1

OUTFILE= 20

OPEN (UNIT=OUTFILE,FILE='L8Mth.DAT',
D.    STATUS='NEW', RECL=250, FORM='FORMATTED')

SKY=22.0  !! !MAG/'^2
EXPMX=1800      ! LMT 16.O SKY -> 8000 E-/PIX-SEC 80K FullWell
! 3 SEC GIVES 50% DUTY CYCLE AT 2X2 BINNING
! EXPMX=30  ! CDT MAUI 22.0 SKY -> 120 e-/PIX-S
INCMX=17  ! 16# OF INC BINS
RANMX=33  ! 33# OF RANGE BINS
RS=850  !!  ! TDI OPTIMUM RANGE GEO39754 (range of desired obj)
BIN=16      !!  ! 37K/20K-2 , 3500-4, 1500-8, 800-16 HOUSTON:>8
DELTR=2  !!  ! GEO/20K-500, 3.5K-50, 1.5K/800-25 400-10
DELT=3  !!  ! INCLINATION INCREMENT 37K=1, all else 2
SITE=0.0    ! SITE LATITUDE  (Use 0.0 for everything)
SI=0.0      ! SEARCH INCLINATION (Inc of desired object)
PI=3.141592
SOLMAG=-26.74

WEARTH=0.004178 !ANG ROT EARTH IN 23:56
GM=398666  !KM3/S2
RE=6378      !KM

```

```

PS=6040      !PIX/DEG  LMT
FW=0.339     !DEG
PXA=0.355    !”2/PIX
READ=28.1    !e-/PIX
DARK=0.0
SN=6.31      !2MAG ABOVE NOISE
PF=2/(3*PI**2) !LAMBERT PF@90DEG !SPHERE=1/4PI (L<Specular)
! PF=2/(PI**2) !LAMBERT PF @ 15DEG
ALB=0.1
! PS=313      !PIX/DEG  CDT
! FW=1.84
! PXA=132
! READ=100
! DARK=50

R=READ
DR=DARK

ALPHA=5.5135E+10 !10,000 P/SEC-CM^2 NM*MIRRORA*BW*THROUGH*QE
! ALPHA=5.702E+8 !CDT
A=ALPHA

DATA(2,1)=0

DO I=1,INCMX
DATA(I+2,1)=I*DELTI
END DO

DO J=1,RANMX
! RO=(RS-DELTR*((RANMX-1)/2))+DELTR*(J-1)
RO=250*(1.191**(J-1))
LRO=LOG10(RO)
DATA(2,J+1)=LRO
END DO

DO I=1,INCMX
DO J=1,RANMX

! RO=(RS-DELTR*((RANMX-1)/2))+DELTR*(J-1)
! RO=500+DELTR*(J-1)
RO=250*(1.191**(J-1))
INC=(I-1)*DELTI*PI/180
RCM=RO*100000

```

```

WST=(180/PI)*(GM/((RE+RS)*RS**2))**0.5
WEWT=(WST*COS(SI))-WEARTH*COS(SITE) ! ew motion of obj'
WNST=WST*SIN(SI)
WS=(WEWT**2+WNST**2)**0.5

```

```

WOT=(180/PI)*(GM/((RE+RO)*RO**2))**0.5
WEW=(WOT*COS(INC))-WEARTH*COS(SITE)
WNS=WOT*SIN(INC)
WO=(WEW**2+WNS**2)**0.5 ! Obj motion wrt telescope

```

```

SIGEXP=FW/(WO*COS(INC))
SIGEXPMX=SIGEXP
EXP=FW/WS
EXPHLD=EXP

```

```

IF (EXP .GT. EXPMX) THEN
    EXP=EXPMX
ENDIF

```

```

IF (WEW .LT. 0) THEN
    WO=-WO
END IF

```

```

EPSILON=(FW*PS/BIN)*(((WS/WO-1)**2+((TAN(INC-SI))**2))**0.5)
NPIX=1+EPSILON

```

```

IF (SIGEXP .GT. EXPMX) THEN
    SIGEXP=EXPMX
    NPIX=1+EPSILON*(EXPMX/SIGEXPMX)
ENDIF

```

```

B=BIN**2
NOISE=((R**2)+(DR*EXP*B)+A*PXA*B*EXP*(10**-0.4*SKY))**0.5
M1=-2.5*LOG10(SN*(NPIX**0.5)*NOISE/(ALPHA*SIGEXP))
M2=M1-1.0 ! Accounts for 6 pix image spread

```

```

DU=RCM*(10**(0.2*(SOLMAG-M2)))
DD=((ALB*PI*PF/4)**0.5)
D=DU/DD

```

```

! DU=RCM*(10**(SOLMAG/5))*(SN**0.5)*(NPIX**0.25)*(NOISE**0.5)
DD=((ALPHA*SIGEXP*ALB*PI*PF/4)**0.5)
DATA(I+2,J+1)=M1

```



```

!          WRITE(*,*)DATA(I,J)
          WRITE(*,*) RO,INC
          WRITE(*,*) EXP, SIGEXP, EXPHLD, BIN
          WRITE(*,*) WO, WS, D, M1, NPIX

          END DO
END DO

!  WRITE(*,*) EXP

INCMX1=INCMX-1
DCHR(1)(1:250)= ' '
DCHR(2)(1:250)= ' '
WRITE(UNIT=DCHR(1)(1:7),FMT='(I6)') INCMX1
WRITE(UNIT=DCHR(1)(8:14),FMT='(I6)') RANMX

DO J=1, RANMX
J1=J-1
WRITE (UNIT=DCHR(2)(1+J1*7:7+J1*7),FMT='(F7.2)') DATA(2,J)
END DO

DO I=3, INCMX

DCHR(I)(1:250)= ' '

DO J=1, RANMX

J1=J-1

WRITE (UNIT=DCHR(I)(1+J1*7:7+J1*7),FMT='(F6.1)')DATA(I,J)

END DO
END DO

DO I=1,INCMX

WRITE(20,200) DCHR(I)
FORMAT(A250)

END DO
END

```

APPENDIX C1

(Original 3 Element Corrector Parameter File)

(Values in parentheses are manufactured values)

3-m POTTER TEL CORRECTOR, Radii from APS Tooling List; 60 mm field.

STOP IS SURF 2

IMAGE IS 11

PXRAY 0.1500000E+04 0.1477440E+17 0.5760000E-02 0.2659392E+02

WVLNS. 0.5000000E+03 0.4000000E+03 0.8000000E+03 0.6000000E+03
0.7000000E+03

1 0 CURVATURE DATA	0.0000000E+00 0.0000000E+00 0.0000000E+00
2 0 AXIAL SEPARATION	-0.2565000E+19
3 0 AIRSPACE	

1 1 CURVATURE DATA	0.0000000E+00 0.0000000E+00 0.0000000E+00
2 1 AXIAL SEPARATION	-0.4617000E+04
3 1 AIRSPACE	

1 2 CURVATURE DATA	0.1112664E-03 0.8987439E+04 -0.1000000E+01
2 2 AXIAL SEPARATION	0.4048655E+04
3 2 REFLECTING SURF.	-0.1000000E+01 -0.1000000E+01 -0.1000000E+01
	-0.1000000E+01 -0.1000000E+01
4 2 TILT ANGLES- 123	0.0000000E+00 0.0000000E+00 0.0000000E+00

1 3 RADIUS DATA	0.5000000E-02 0.2000000E+03 0.0000000E+00 (0.19967E+3)
2 3 AXIAL SEPARATION	0.2400000E+02 (0.2422E+02)
3 3 INDICES	-0.1521415E+01 -0.1530852E+01 -0.1510779E+01 -0.1516294E+01 -0.1513065E+01
3 3 ABBE NUMBER = 64.13 Glass type BK7	(64.2)

1 4 RADIUS DATA	0.4945598E-02 0.2022000E+03 0.0000000E+00 (0.202575E+03)
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2 4 AXIAL SEPARATION	0.2408604E+03		
3 4 AIRSPACE			
1 5 RADIUS DATA	0.2689618E-02	0.3718000E+03	0.0000000E+00
		(0.37192E+03)	
2 5 AXIAL SEPARATION	0.3920000E+02		
	(0.39115E+02)		
3 5 INDICES	-0.1521415E+01	-0.1530852E+01	-0.1510779E+01
	-0.1516294E+01	-0.1513065E+01	
3 5 ABBE NUMBER =	64.13	Glass type BK7	(64.2)
1 6 RADIUS DATA	0.1171783E-01	0.8534000E+02	0.0000000E+00
		(0.8541E+02)	
2 6 AXIAL SEPARATION	0.1004035E+03		
3 6 AIRSPACE			
1 7 RADIUS DATA	0.6995453E-02	0.1429500E+03	0.0000000E+00
		(0.143034E+03)	
2 7 AXIAL SEPARATION	0.1570000E+02		
	(0.1576E+02)		
3 7 INDICES	-0.1521415E+01	-0.1530852E+01	-0.1510779E+01
	-0.1516294E+01	-0.1513065E+01	
3 7 ABBE NUMBER =	64.13	Glass type BK7	(64.2)
1 8 RADIUS DATA	-0.2834467E-02	-0.3528000E+03	0.0000000E+00
		(0.3527E+03)	
2 8 AXIAL SEPARATION	0.3680925E+02		
3 8 AIRSPACE			
1 9 CURVATURE DATA	0.0000000E+00	0.0000000E+00	0.0000000E+00
2 9 AXIAL SEPARATION	0.3175000E+01		
3 9 INDICES	-0.1462319E+01	-0.1470114E+01	-0.1453310E+01
	-0.1458035E+01	-0.1455289E+01	
3 9 ABBE NUMBER =	67.86	Glass type FS	(67.8)
1 10 CURVATURE DATA	0.0000000E+00	0.0000000E+00	0.0000000E+00
2 10 AXIAL SEPARATION	0.9525000E+01		
3 10 AIRSPACE			
1 11 CURVATURE DATA	0.0000000E+00	0.0000000E+00	0.0000000E+00

potopt500.out

centroid for all rays : 0.0000 -0.0011

colours : 400.00 500.00 650.00 800.00 1000.00

y offsets for each height, colour

D.1	0.0110	0.0101	0.0082	0.0067	0.0050
18.49	-0.0056	-0.0045	-0.0044	-0.0046	-0.0050
15.05	-0.0061	-0.0049	-0.0047	-0.0048	-0.0050
7.50	-0.0040	-0.0033	-0.0031	-0.0031	-0.0031
D.1	0.0000	0.0000	0.0000	0.0000	0.0000

Maximum spot diameter = 0.14419431468227

Maximum spot sizes ...

D.1	0.0636	0.0827	0.0914	0.0980
D.1	0.1258	0.1176	0.1150	0.1140
D.1	0.1073	0.1030	0.1025	0.1033
D.1	0.0494	0.0528	0.0565	0.0608
D.1	0.0061	0.0105	0.0180	0.0257

80 percent-energy spot sizes ...

D.1	0.0333	0.0447	0.0512	0.0576
D.1	0.0292	0.0236	0.0229	0.0246
D.1	0.0257	0.0189	0.0182	0.0202
D.1	0.0143	0.0077	0.0090	0.0113
D.1	0.0047	0.0055	0.0068	0.0080

Half-energy spot sizes ...

D.1	0.0211	0.0252	0.0283	0.0318
D.1	0.0153	0.0103	0.0110	0.0129
D.1	0.0136	0.0093	0.0095	0.0114
D.1	0.0078	0.0056	0.0046	0.0061
D.1	0.0037	0.0036	0.0052	0.0068

total number of rays = 9646

APPENDIX C2

(Modified 4 Element Corrector Parameter File)

07/11/97

NASA-LMT

Modification to the existing 3 Element Wynne Corrector

Elements 1 and 2 retained

Elements 3 and 4 new

System saved in file CONBK.LEN(5)

CODE V> lis

LMT, 4-el, 3rd aspheric

	RDY	THI	RMD	GLA	CCY	THC	GLC
OBJ:	INFINITY	INFINITY			100	100	
STO:	-8988.14679	-4045.616841	REFL			0	100

CON:

K : -1.000000 KC : 100

2:	-199.67000	-24.220000	BK7_SCHOTT (old)	100	100
3:	-202.57500	-240.860400	AIR	100	100
4:	-371.92000	-39.115000	BK7_SCHOTT (old)	100	100
5:	-85.41000	-66.694774	AIR	100	0
6:	-77.86514	-10.000000	BK7_SCHOTT (new)	0	0
7:	-66.04087	-33.063174	AIR	0	0

CON:

K : -0.371270 KC : 0

8:	-70.08598	-18.518725	BK7_SCHOTT (new)	0	0
9:	-313.45086	-18.422555	AIR	0	0
10:	INFINITY	-7.000000	'gg455'	100	100
11:	INFINITY	-5.000000		100	100
12:	INFINITY	-3.175000	SILICA_SPECIAL	100	100
13:	INFINITY	-9.525000		100	100
IMG:	INFINITY	0.000000		100	100

SPECIFICATION DATA

EPD	2999.99986					
DIM	MM					
WL	1000.00	800.00	600.00	500.00	400.00	
REF	3					
WTW	1	1	1	1	1	
INI	her					
XAN	0.00000	0.00000	0.00000	0.00000	0.00000	0.00000
	0.00000					
YAN	0.00000	0.08000	0.16000	0.22400	0.32000	
	-0.32000					

APERTURE DATA/EDGE DEFINITIONS

CA	
CIR S1	1500.000000
CIR S2	148.000000
CIR S4	70.000000
CIR S6	53.500000
CIR S7	53.500000
CIR S8	48.500000
CIR S9	48.500000
CIR S2 EDG	148.000000
CIR S4 EDG	70.000000
CIR S6 EDG	53.500000+-0.1 mm
CIR S7 EDG	53.500000+-0.1 mm w/ 4.5mm Land
CIR S8 EDG	48.500000+-0.1 mm
CIR S9 EDG	48.500000+-0.1 mm w/ 3.5mm Land (if necessary)

PRIVATE CATALOG

PWL	1014.00	587.60	480.00
'gg455'	1.530000	1.540000	1.550000

REFRACTIVE INDICES

GLASS CODE	1000.00	800.00	600.00	500.00	400.00
BK7_SCHOTT	1.507502	1.510776	1.516295	1.521414	1.530849
SILICA_SPECIAL	1.450416	1.453318	1.458041	1.462331	1.470122
'gg455'	1.530130	1.532831	1.539300	1.547361	1.569763

INFINITE CONJUGATES

EFL -5152.9216

BFL -9.5684

FFL 50653.7178

FNO 1.7176

IMG DIS -9.5250

OAL -4511.6865

PARAXIAL IMAGE

HT 28.7796

ANG 0.3200

ENTRANCE PUPIL

DIA 2999.9999

THI 0.0000

EXIT PUPIL

DIA 305.1852

THI -533.7669

APPENDIX D

Galaxy Distribution/Clustering for a R.A. and Declination Wedge to $z = 0.033$

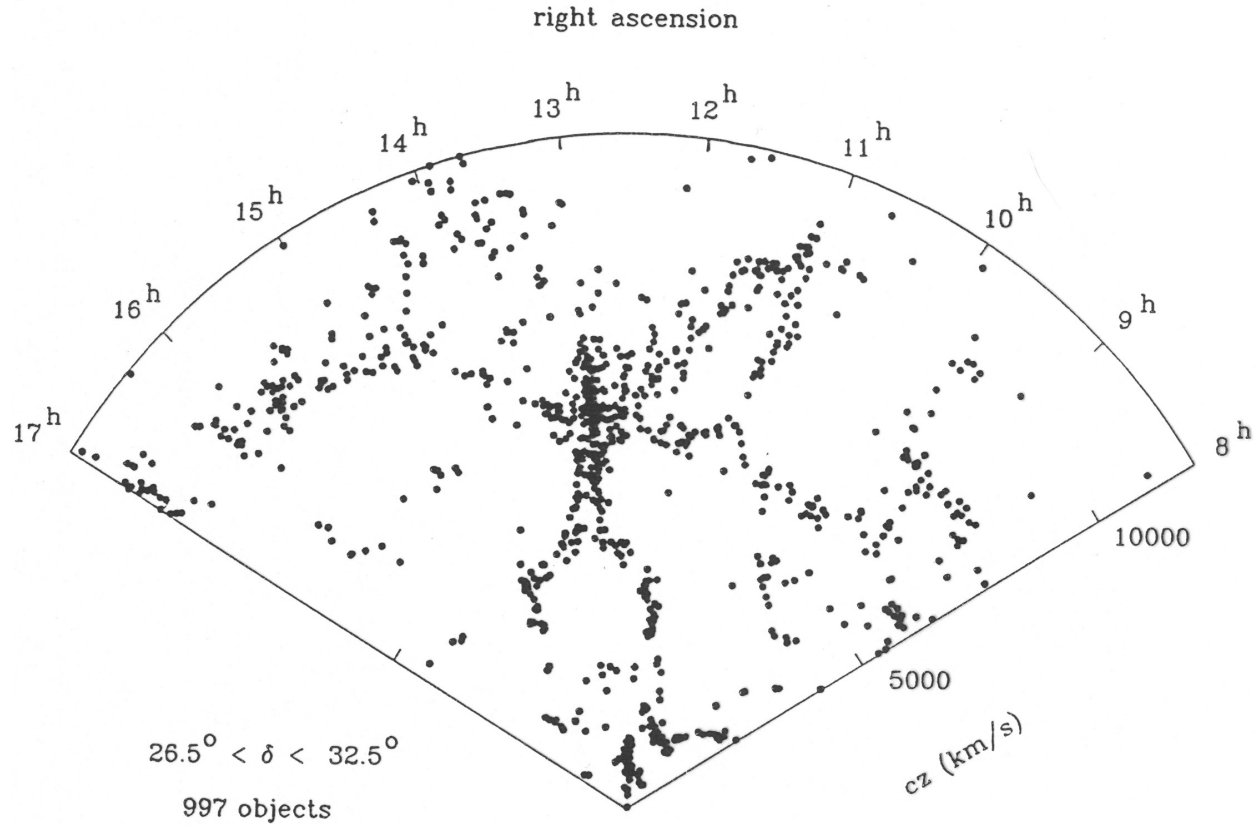


Figure AD-1. A two dimensional representation of a R.A. and Declination wedge showing the distribution of galaxies to $z = 0.033$. Galaxy clustering is clearly evident. The UBC/NASA-LMT multi-narrowband survey will generate a similar figure, but extending 15 times deeper in z (0.5), over 12-18 hr R.A. and ± 16 arcminute declination range centered on $+32.979$ degrees. (Figure credit: de Lapparent et al. 1986).

APPENDIX E

B Band Filter Transmission

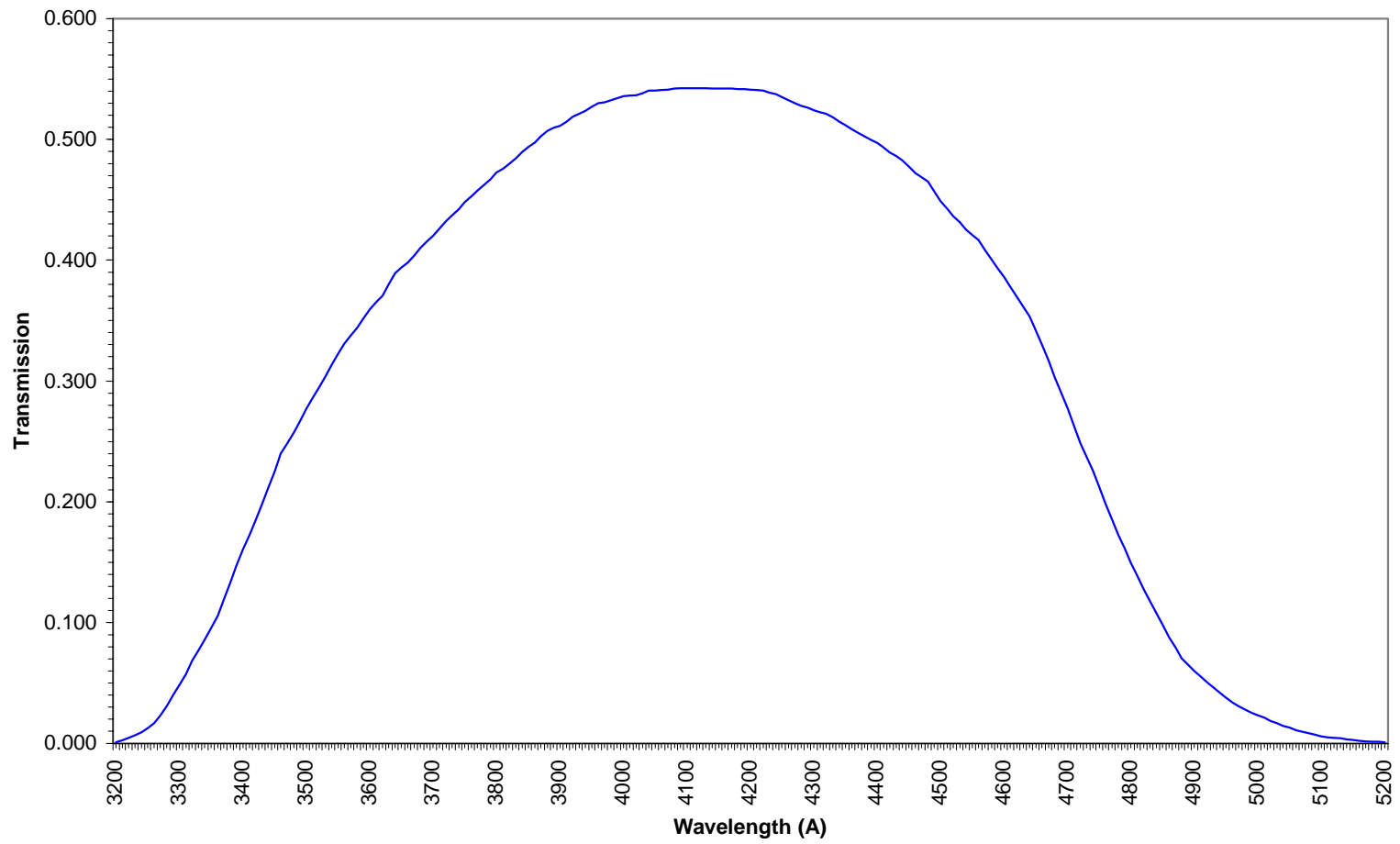


Figure AE-1. Measured broadband transmission for the UBC/NASA B Band filter.

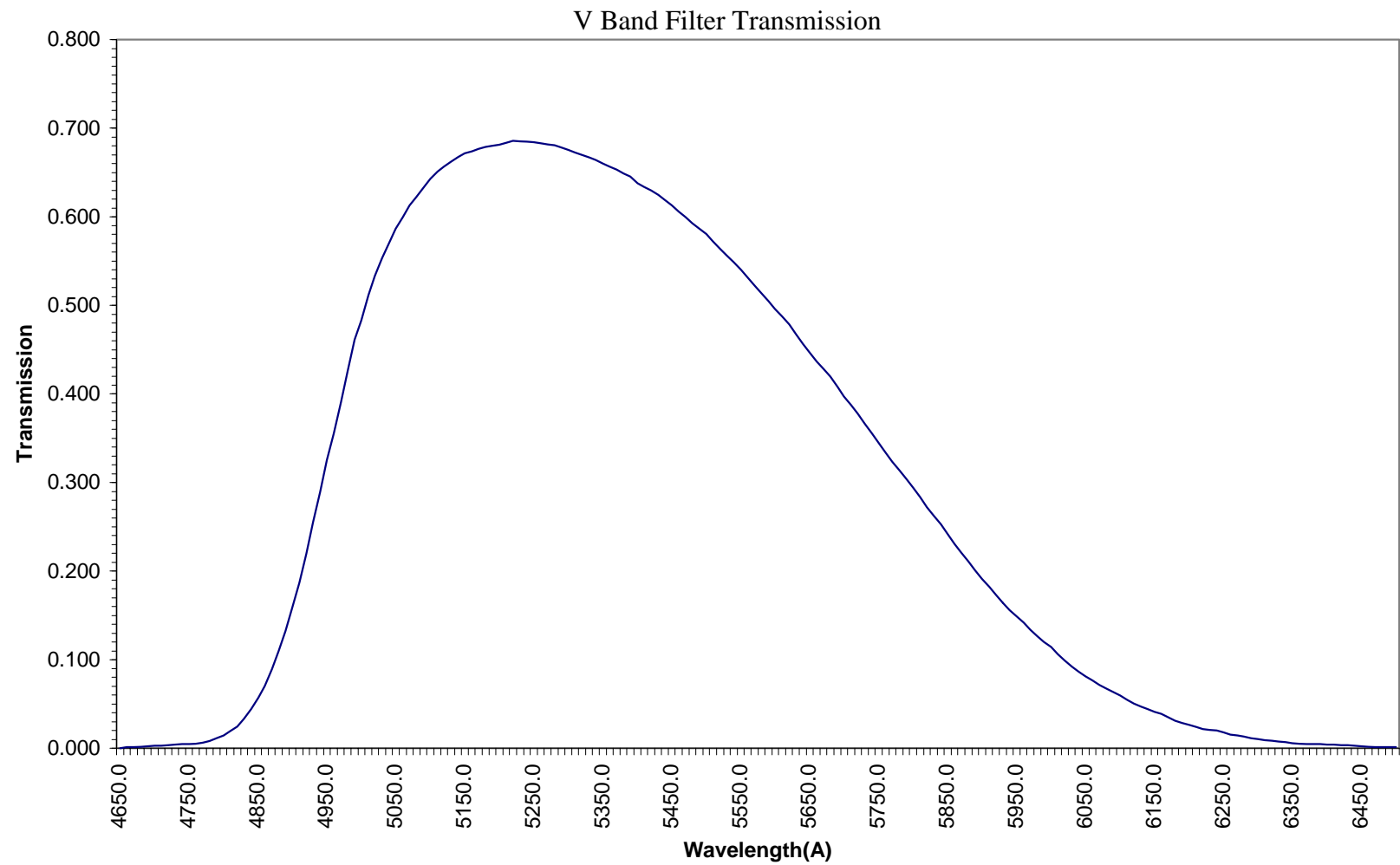


Figure AE-2. Measured broadband transmission for the UBC/NASA V Band filter.

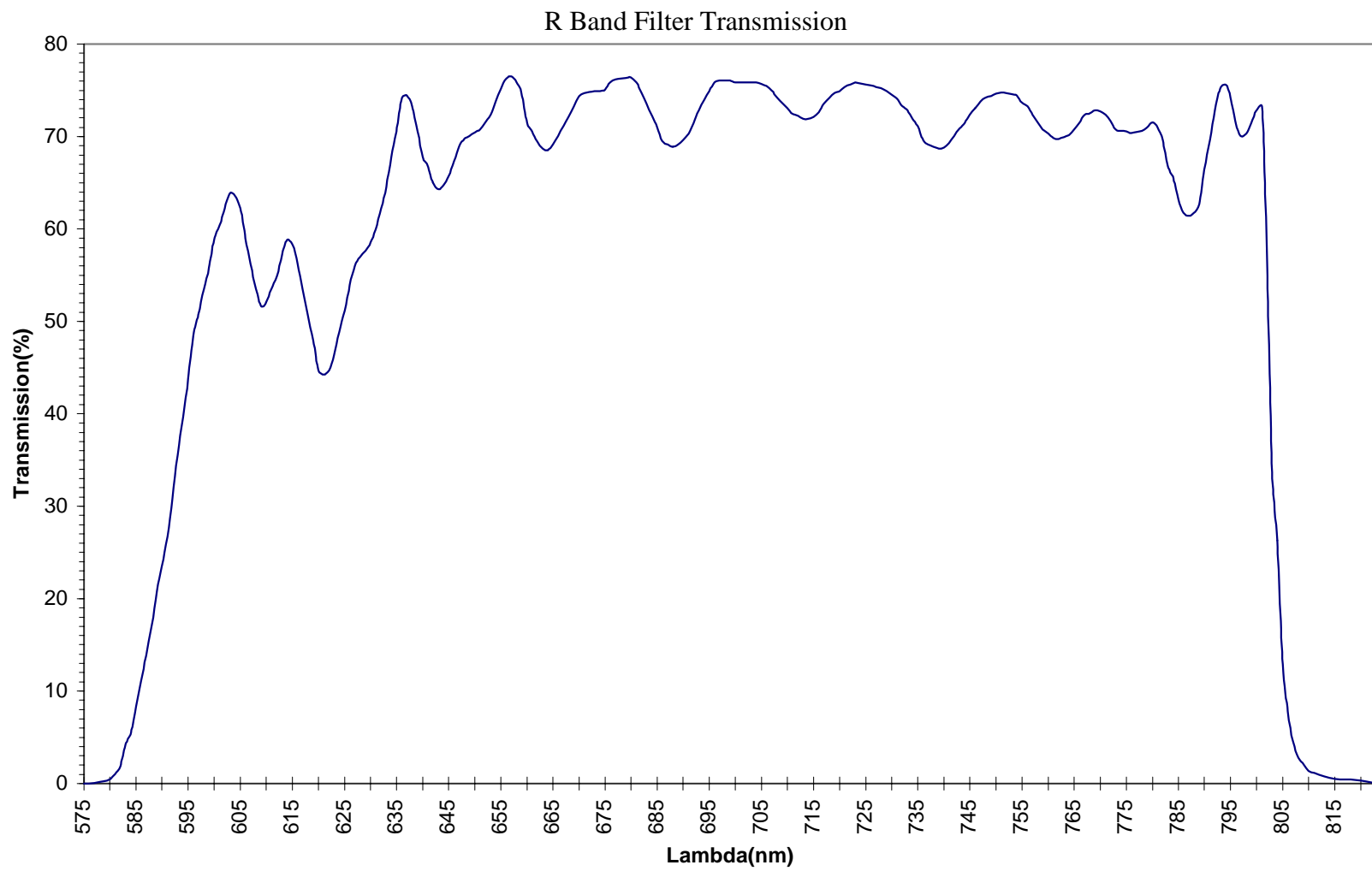


Figure AE-3. Measured broadband transmission for the UBC/NASA R Band filter.

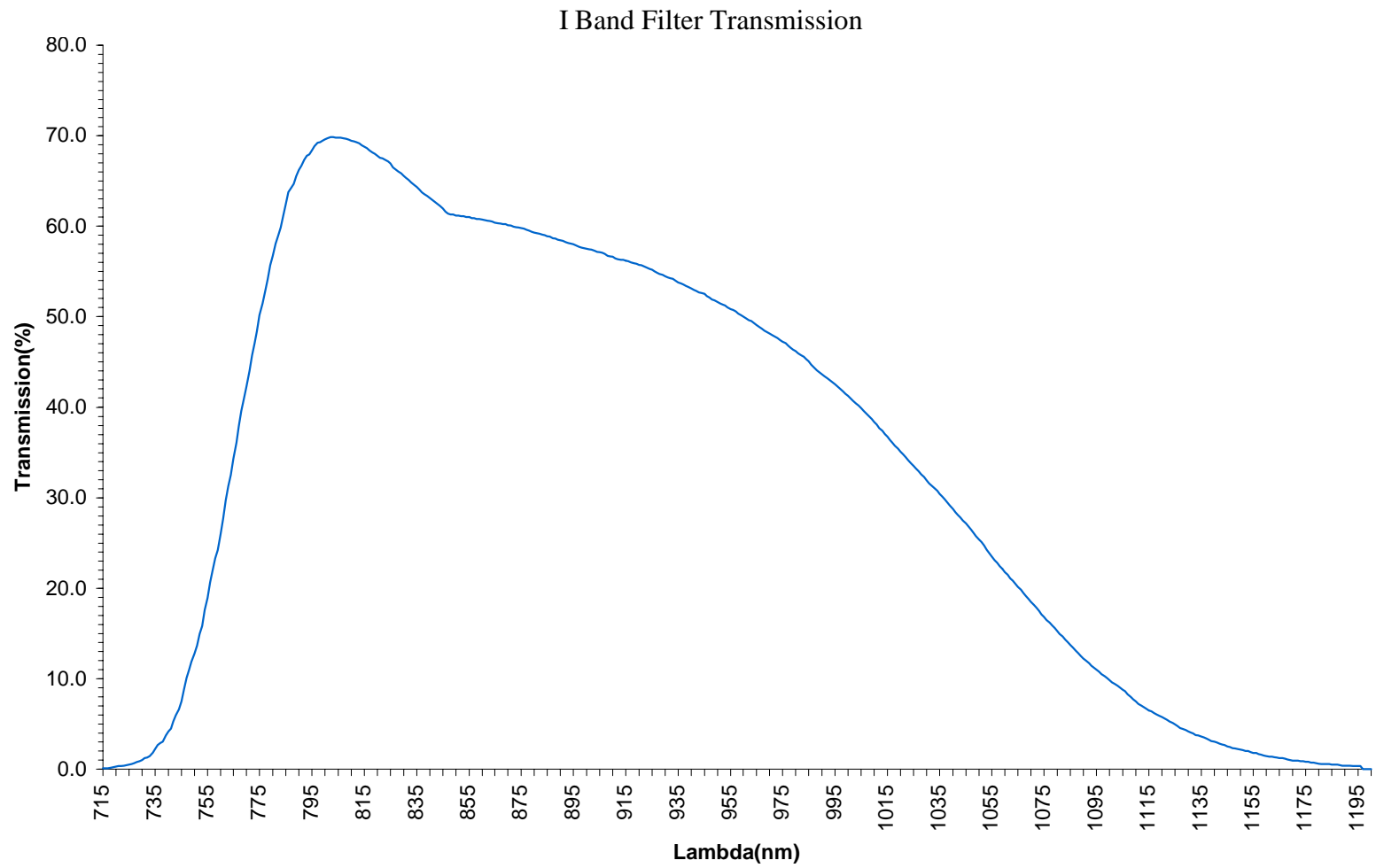


Figure AE-4. Measured broadband transmission for the UBC/NASA I Band filter.

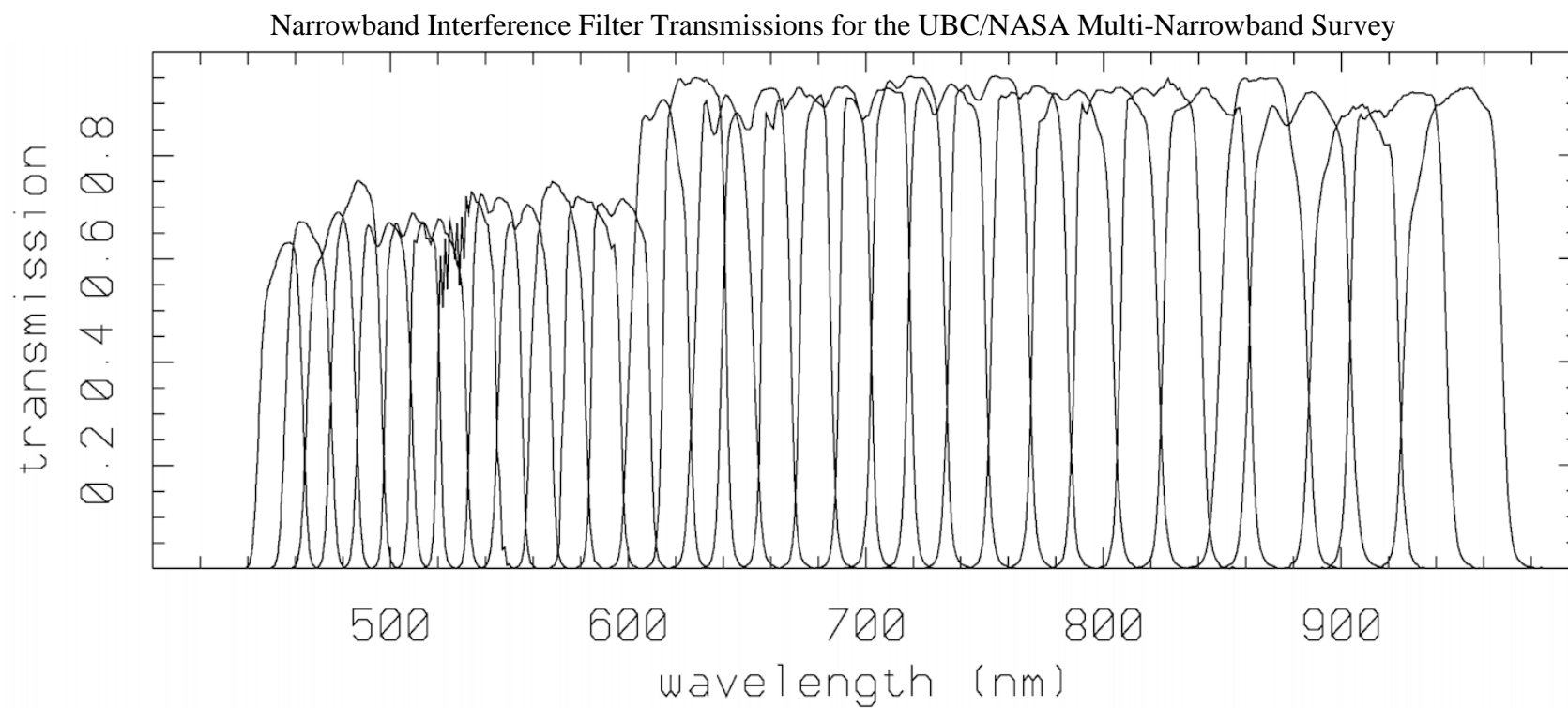


Figure AE-5. Transmissions for the 40 narrowband interference filters used in the UBC/NASA Multi-Narrowband Survey. The filters are spaced equally in log frequency and have equal energy bandpasses at all frequencies. Used together they can yield photometric redshifts (z) with ± 0.02 accuracy.